

Carrier Density Control of Magnetism and Hall Effects in EuTiO_3 Films

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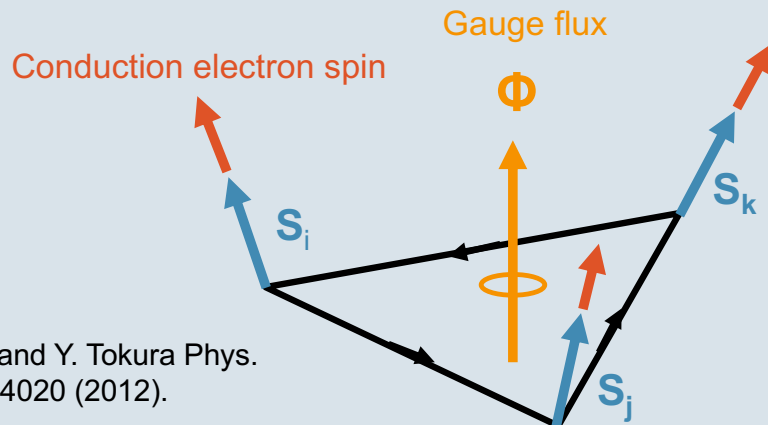
Correlated Electron Systems – Novel Developments Workshop
Fine Theoretical Physics Institute, University of Minnesota
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Outline

- ❑ Spontaneous Hall effects in magnetic materials
- ❑ Rare-earth doped EuTiO_3 films
 - ❑ Fermi surface topology, magnetism, and quantum criticality
- ❑ $\text{SmTiO}_3/\text{EuTiO}_3$ interfaces
- ❑ Summary

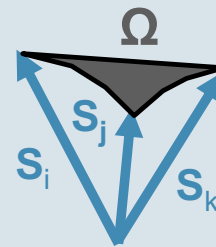
Berry Curvature and Hall Effects

Berry phase associated with a non-coplanar spin structure

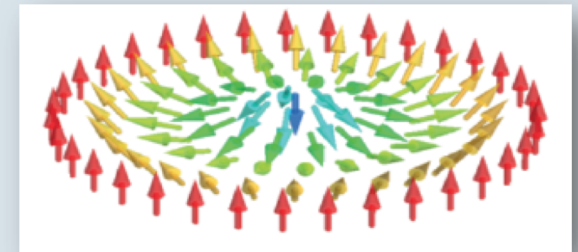


N. Nagaosa and Y. Tokura Phys. Scr. **T146** 014020 (2012).

Spin chirality
 $S_i \cdot (S_j \times S_k)$



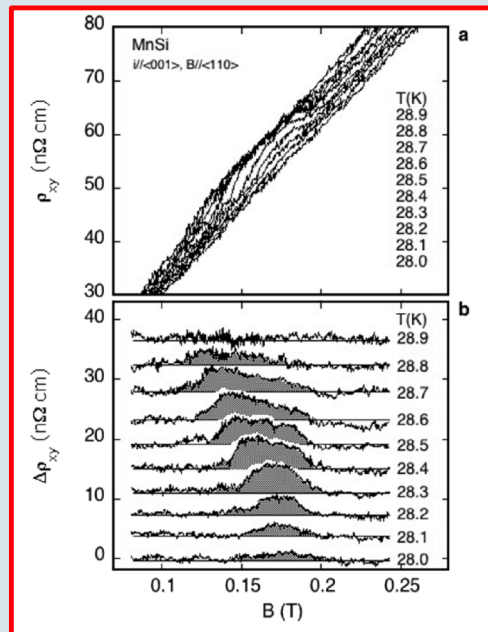
Skyrmion



- ❑ Berry phase of the magnetization texture
- ❑ Electron spins of the conduction electrons are coupled adiabatically and therefore accumulate a Berry phase
- ❑ Effective Lorentz force results in a “Topological Hall effect” P. Bruno, et al., PRL 93, 096806 (2004).

Topological Hall effect

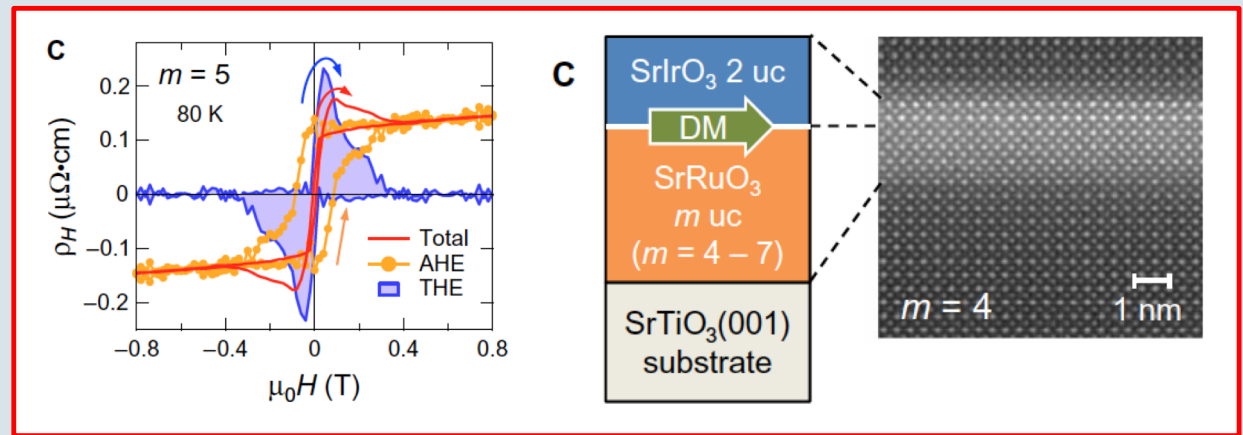
Experimental Observation of the Topological Hall Effect



MnSi - Skymion Phase

A. Neubauer, et al., PRL 102, 186602 (2009).

Theory: Ye et al., PRL 83, 3737 (1999); Y. Taguchi et al., Science 291, 2573 (2001); Onoda et al., PRL 90, 196602 (2003); P. Bruno, et al., PRL 93, 096806 (2004); B. Binz et al., Physica B 403, 1336 (2008).



SrIrO₃/SrRuO₃ Interface Skyrmions

Matsuno et al., Sci. Adv. 2, e1600304 (2016).

Topological Hall effect

Berry Curvature and Hall Effects

Berry phase in reciprocal space

Equation of motion in an external field:

velocity in state \mathbf{k}

$$v_n(\mathbf{k}) = \underbrace{\frac{\delta \epsilon_n(\mathbf{k})}{\hbar \delta(\mathbf{k})}}_{\text{Band dispersion}} - \frac{e}{\hbar} \mathbf{E} \times \underbrace{\Omega_n(\mathbf{k})}_{\text{Anomalous velocity}} \rightarrow \text{Hall effect}$$

R. Karplus and J. M. Luttinger: Phys. Rev. 95 1154 (1954).

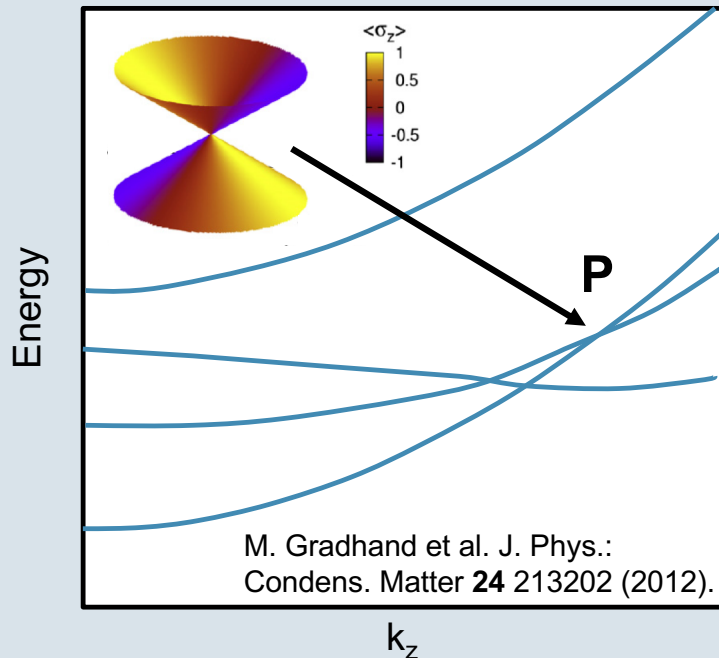
M.-C. Chang and Q. Niu, PRL 75, 1348 (1995).

F. D. M. Haldane, PRL 93, 206602 (2004).
Nagaosa et al., Rev. Mod. Phys. 82, 1539 (2010).

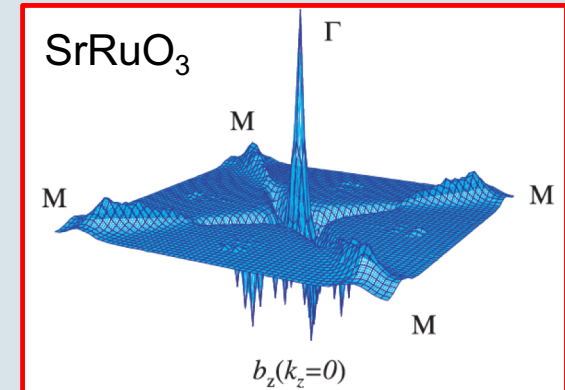
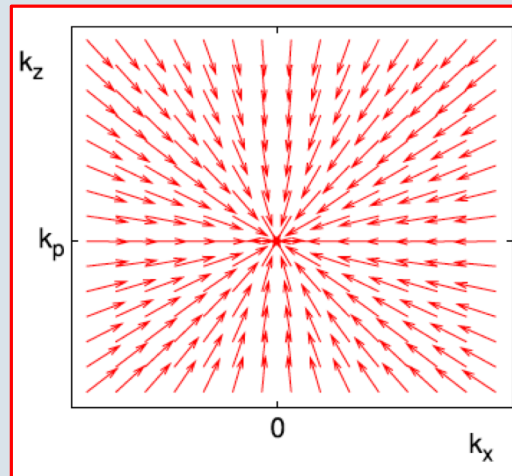
- ❑ $\Omega_n(\mathbf{k})$ is the Berry curvature of the n^{th} band
- ❑ Acts in reciprocal space like a magnetic flux density in real space: “reciprocal space magnetic field”
- ❑ $\Omega_n(\mathbf{k})$ is a band structure property, in addition to the usual band dispersion
- ❑ Non-zero in systems with broken time-reversal symmetry (magnetic materials) OR broken space inversion symmetry

Intrinsic Anomalous Hall Effect

Berry Curvature in Reciprocal Space



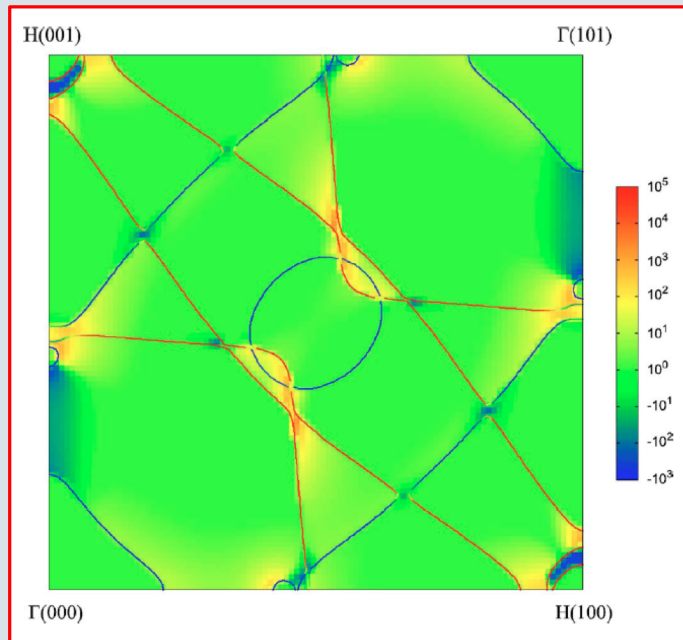
Magnetic monopoles in reciprocal space



Fang et al., Science 302, 92 (2003).

- ❑ Band crossings or avoided band crossings cause large variations in spin with \mathbf{k} and are sources/sinks of Berry curvature
- ❑ Magnetic monopoles in reciprocal space

Berry Curvature in Reciprocal Space



The calculated Berry curvature is shown in Fig. 6. It can be seen that the regions in which the Berry curvature is small (light green regions) fill most of the plane. The largest values occur at the places where two Fermi lines approach one another, consistent with the discussion of Fig. 3. Of special importance are the avoided crossings between two bands having the same sign of spin, or between two bands of opposite spin. Examples of both kinds are visible in the figure, and both tend to give rise to very large contributions in the region of the avoided crossing. Essentially, the spin-orbit interaction causes the character of these bands to change extremely rapidly with \mathbf{k} near the avoided crossing; this is the origin of the large Berry curvature.

Wang et al., PRB **74**, 195118 (2006).

- ❑ Band crossings or avoided band crossings cause large variations in spin with \mathbf{k} and are sources/sinks of Berry curvature

Anomalous Hall Effect

$$v_n(\mathbf{k}) = \underbrace{\frac{\delta\epsilon_n(\mathbf{k})}{\hbar\delta(\mathbf{k})}}_{\text{Band dispersion}} - \frac{e}{\hbar} \mathbf{E} \times \underbrace{\Omega_n(\mathbf{k})}_{\text{Anomalous velocity}} \rightarrow \text{Anomalous Hall effect}$$

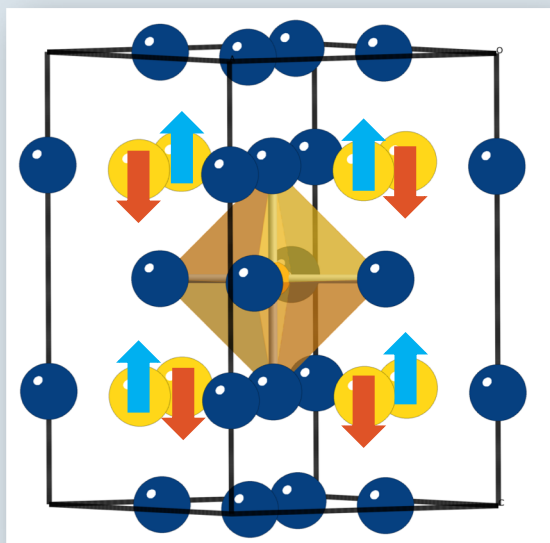
- ❑ $\Omega_n(\mathbf{k})$ is a band structure property, in addition to the usual band dispersion
- ❑ Anomalous Hall effect even in antiferromagnets with small net magnetization

H. Chen, Q. Niu, and A. H. MacDonald, Phys. Rev. Lett. 112, 017205 (2014); J. Kübler and C. Felser, Europhys. Lett. 108, 67001 (2014). S. Nakatsuji, N. Kiyohara and T. Higo, Nature 527, 212 (2015).

- ❑ Expect the intrinsic anomalous Hall effect to depend on Fermi level position
Z. Fang et al., Science 302, 92 (2003).

→ Control via carrier density (doping, electric field effect, ...)

Rare Earth Doped EuTiO_3



Eu 4f spins ($S = 7/2$)

G-type Antiferromagnetic order: $T_N = 5.5 \text{ K}$

Cubic at room temperature

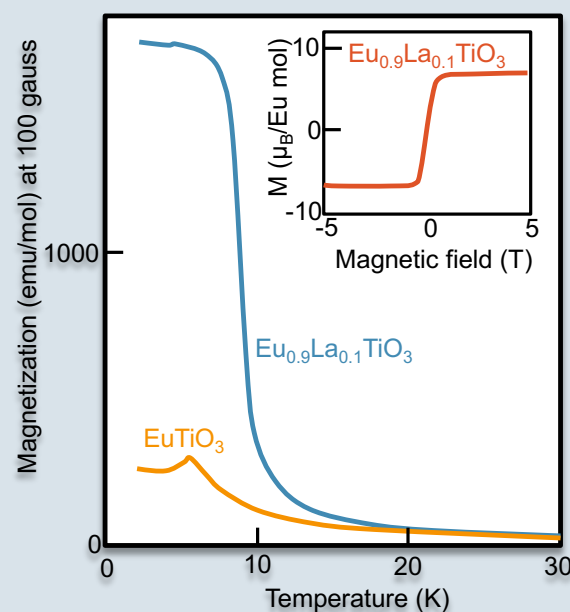
Antiferrodistortive transition ($I4/mcm$)

PHYSICAL REVIEW B

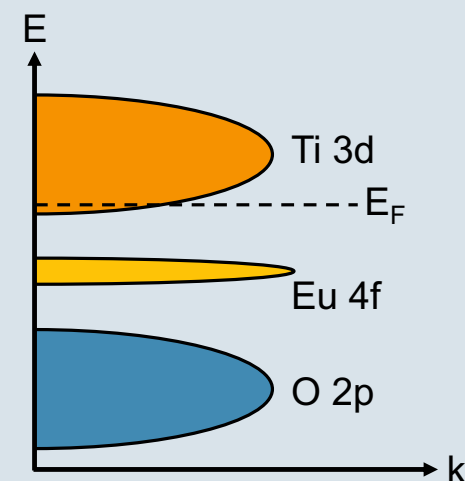
VOLUME 60, NUMBER 22

1 DECEMBER 1999-II

Transport and magnetic properties of a ferromagnetic metal: $\text{Eu}_{1-x}\text{R}_x\text{TiO}_3$



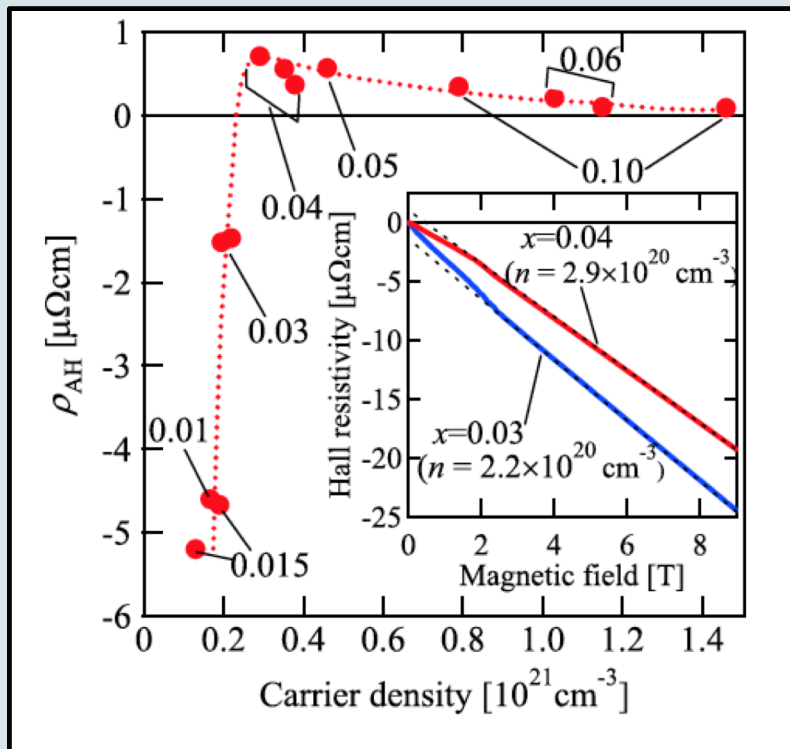
T. Katsufuji and Y. Tokura, Phys. Rev. B 60, R15021 (1999).



Carrier mediated (ferro)magnetism and transport?

R:EuTiO_3 ($\text{R} = \text{La, Gd}, \dots$)

Rare Earth Doped EuTiO_3



K. S. Takahashi et al., PRL 103, 057204 (2009).

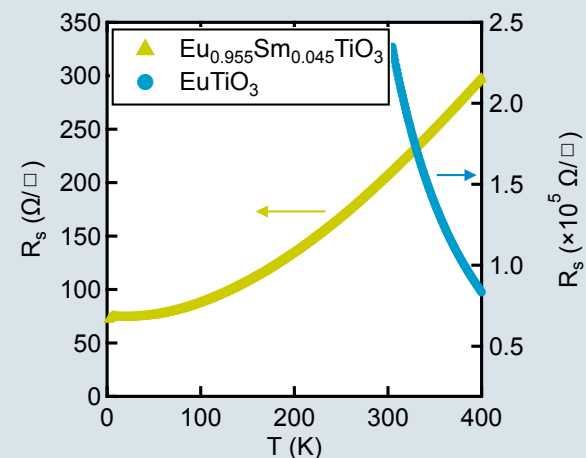
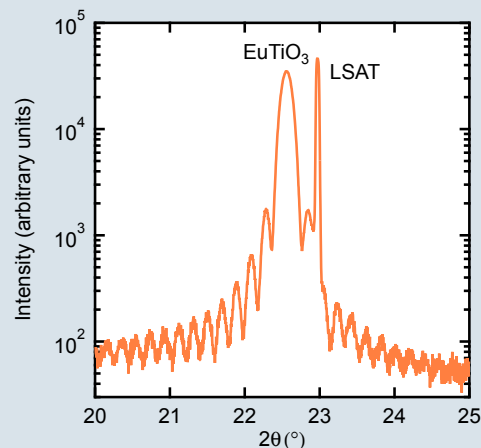
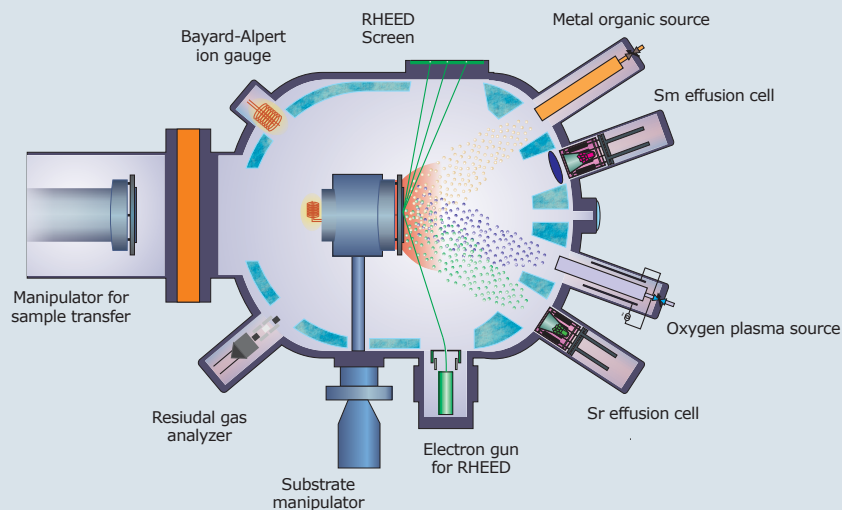
- ❑ EuTiO_3 : Insulating AFM below 5 K
- ❑ $\text{Eu}_{1-x}\text{La}_x\text{TiO}_3$
- ❑ Metallic (anti-)ferromagnet
T. Katsufuji and Y. Tokura, Phys. Rev. B 60, R15021 (1999).
- ❑ Sign change in anomalous Hall effect with carrier density
- ❑ Suggests intrinsic origin of AHE

Why Rare Earth Doped EuTiO_3 ?

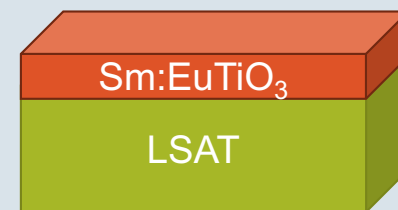
- ❑ $\text{Eu}_{1-x}\text{R}_x\text{TiO}_3$ (R = trivalent rare earth ion such as Sm or La)
- ❑ Metallic (anti-)ferromagnet T. Katsufuji and Y. Tokura, Phys. Rev. B 60, R15021 (1999).

- ❖ Sign change in AHE with carrier density suggests intrinsic, topological origin
 - **Band crossing**
 - Possibly other types of interesting phenomena: changes in the Fermi surface topology by doping, magnetic field, ...
- ❖ Low carrier densities (compared to typical metals)
 - Low Fermi level
 - Magnetic energy scales can drive transitions
- ❖ Low carrier densities → electric field-effect control

EuTiO₃ Films by Molecular Beam Epitaxy

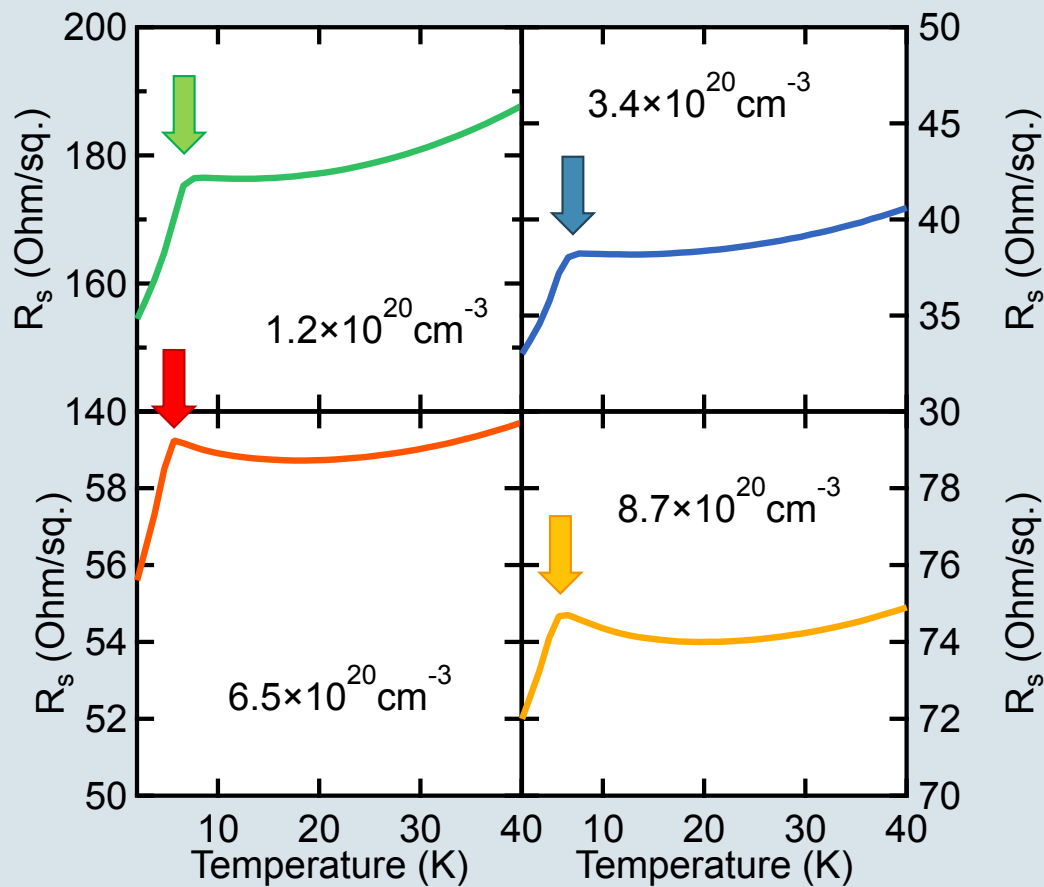


- ❑ High-quality epitaxial films using a metal-organic precursor (TTIP) as Ti and oxygen source
- ❑ Electron doping with Sm causes films to become metallic
- ❑ Vary doping concentration: $1 - 9 \times 10^{20} \text{cm}^{-3}$

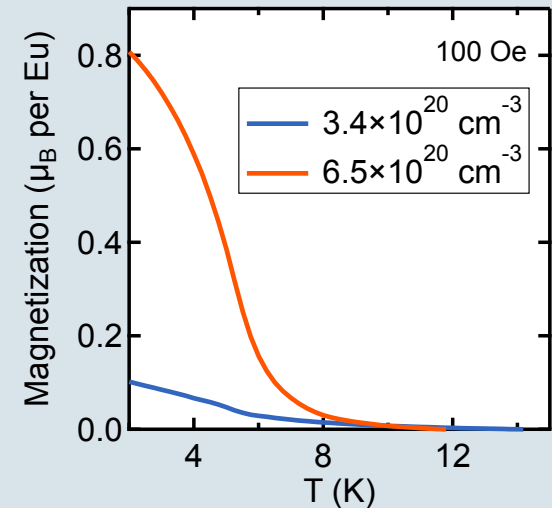


K. Ahadi, L. Galletti, and S. Stemmer, Appl. Phys. Lett. 111, 172403 (2017).

Carrier Mediated Magnetism?

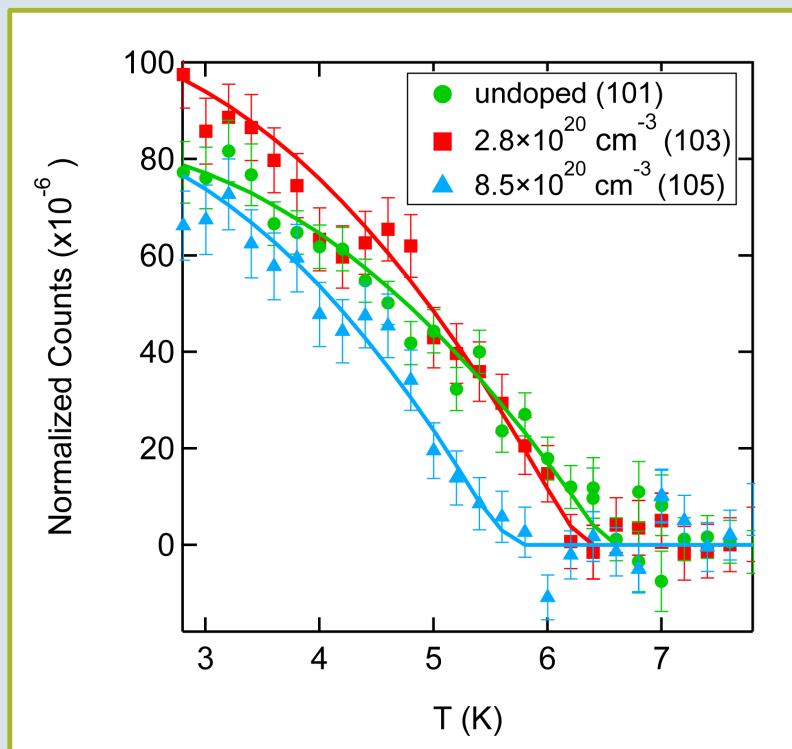


- ☐ Resistance anomaly at the magnetic ordering temperature
- ☐ Magnetic ordering temperature barely changes with doping concentration
- ☐ Similar to T_N of undoped EuTiO_3
- ☐ Change in magnetic order with doping?



K. Ahadi, et al., APL Materials 6, 056105 (2018).

Carrier Mediated Magnetism?



- ❑ Neutron diffraction measurements as a function of doping
- ❑ Magnetic Bragg peaks at tetragonal [odd, 0, odd] (pseudocubic $[\frac{1}{2} \frac{1}{2} \frac{1}{2}]$ -type) observed in all samples below T_N
- ❑ Persistent antiferromagnetic order with carrier substitution
- ❑ Only small change in T_N with doping

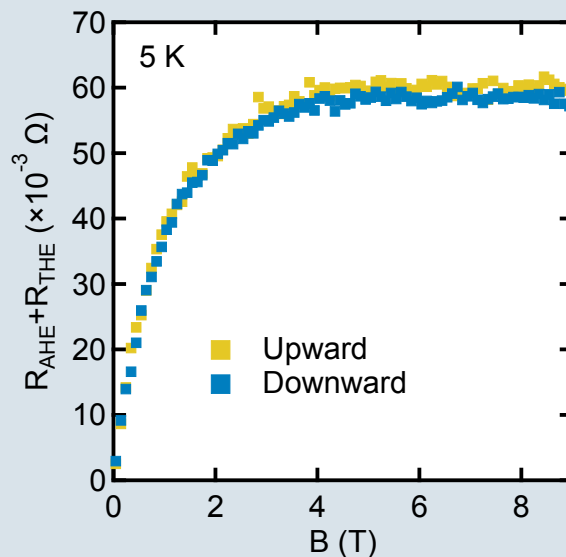
K. Ahadi, et al., APL Materials 6, 056105 (2018).

Spontaneous Hall Effects

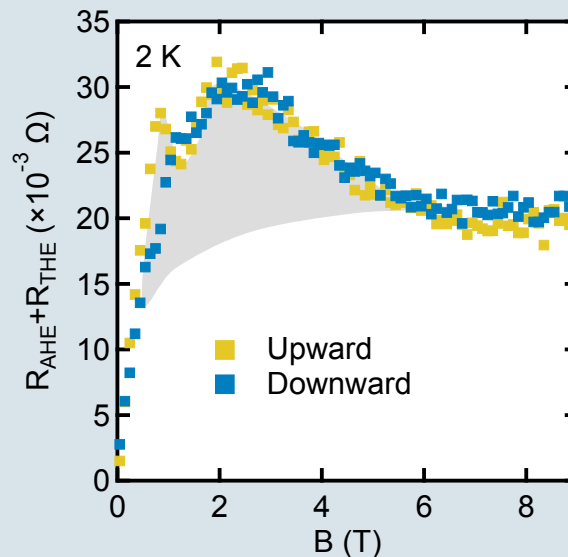
$$R_{xy} = R_0 B + R_{\text{AHE}} + R_{\text{THE}}$$

$$n = 9 \times 10^{20} \text{cm}^{-3}$$

AHE

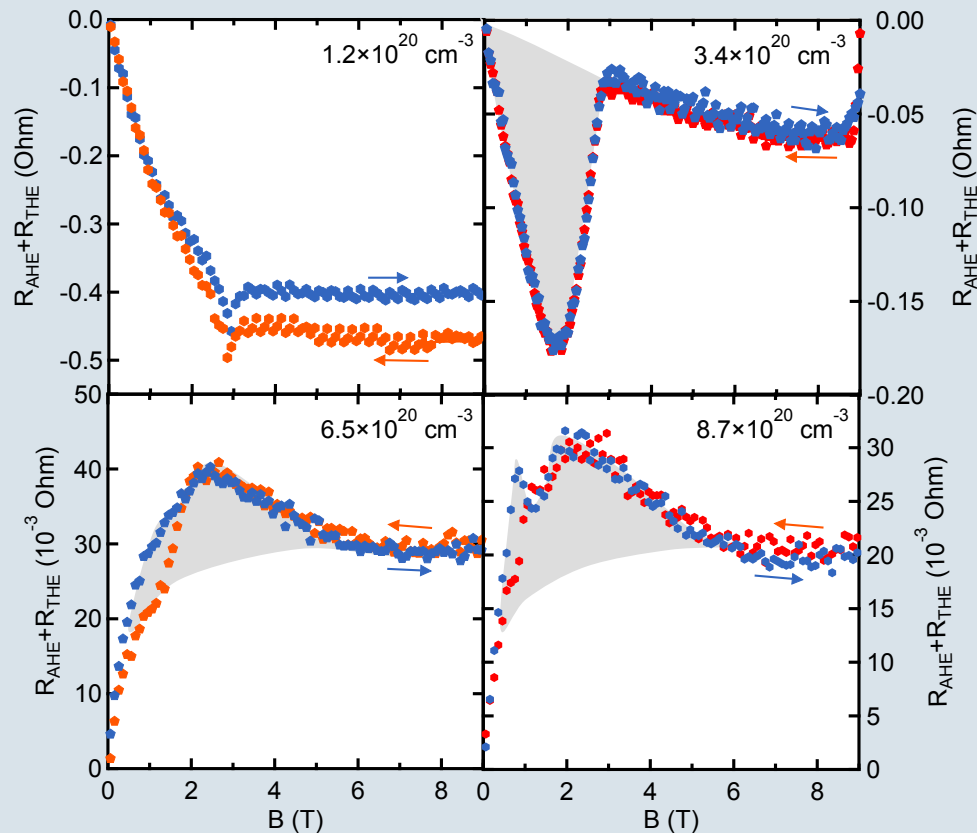


AHE+THE



- Below ~ 30 K: Anomalous Hall effect (AHE)
- Below 5 K: Additional non-monotonic contribution

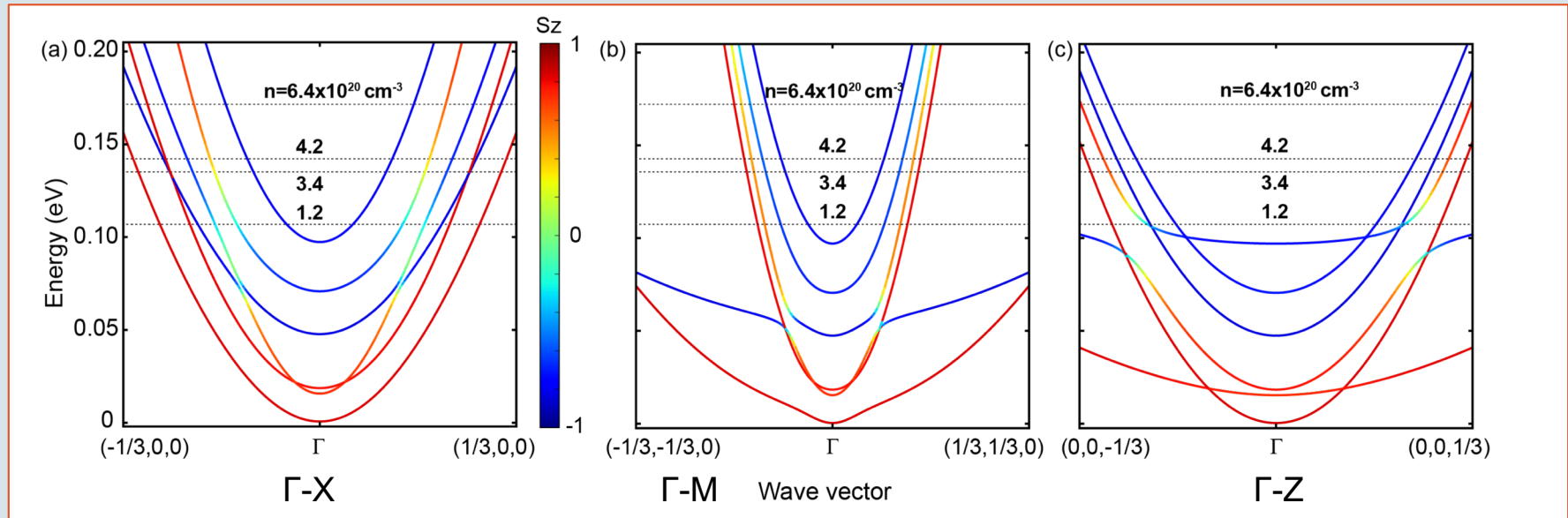
Spontaneous Hall Effects



K. Ahadi, et al., APL Materials 6, 056105 (2018).

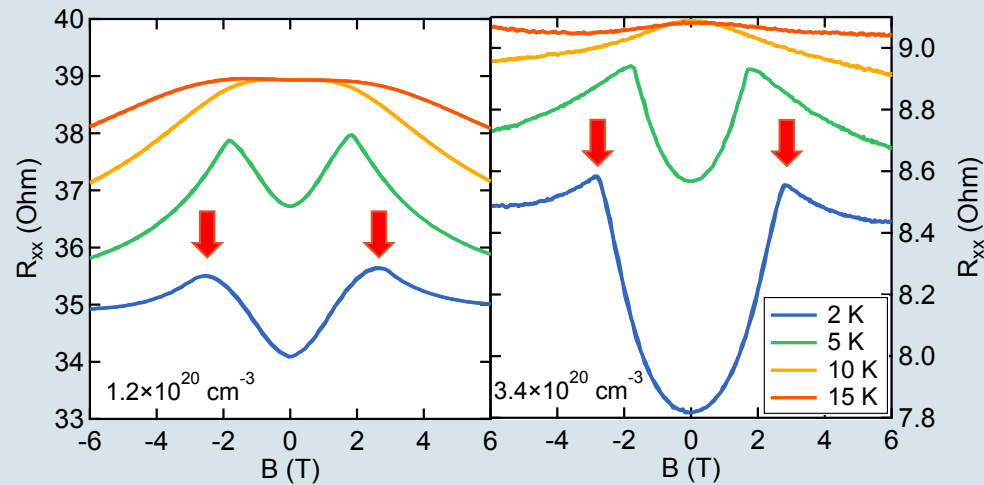
- ☐ Different carrier densities
- ☐ Both spontaneous Hall effects at all carrier densities
- ☐ Despite small net magnetization
- ☐ Sign change in BOTH effects at $\sim 4 \times 10^{20} \text{ cm}^{-3}$
- ☐ **Change in both reciprocal (and real space) Berry curvatures**
- ☐ Reciprocal space: Weyl points?

Anomalous and Topological Hall Effects



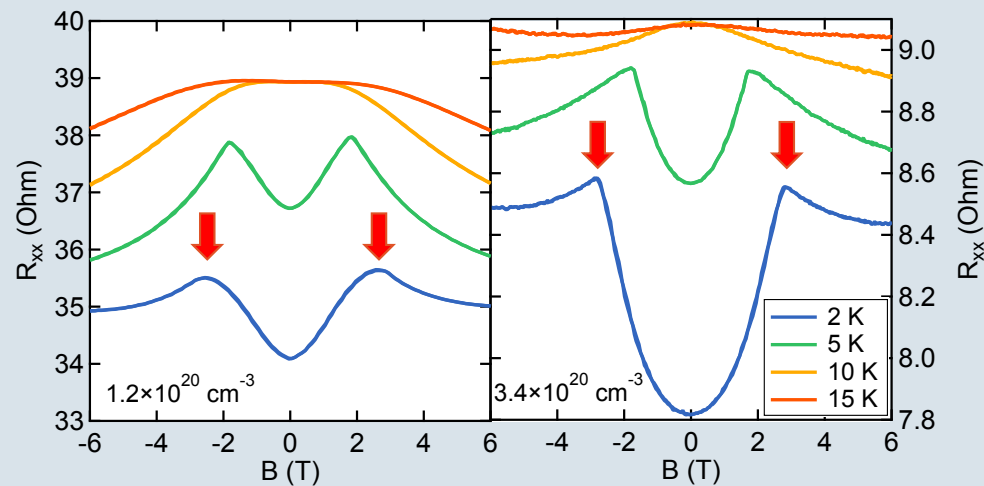
- ❑ DFT calculations (Janotti group)
- ❑ Ti-3d conduction band with spin-up crosses another Ti-3d conduction-band with spin-down along the Γ -X (in-plane) direction, in a spin-momentum locked configuration at the Fermi level for $\sim 4 \times 10^{20} \text{ cm}^{-3}$
- ❑ **Source of reciprocal space Berry curvature** K. Ahadi, et al., APL Materials 6, 056105 (2018).

Magnetoresistance

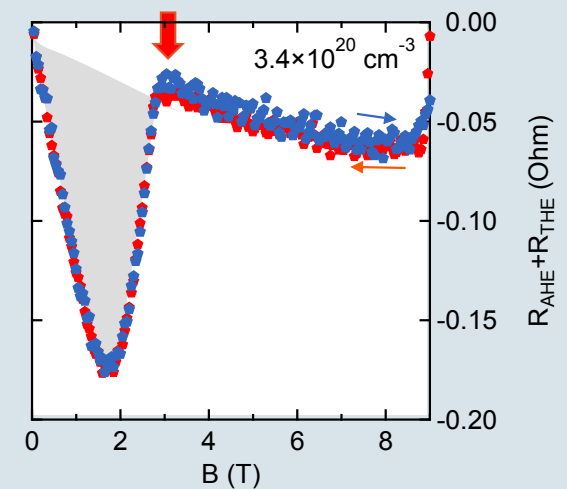


- ☐ Low doping
- ☐ Low-magnetization state
- ☐ Metamagnetic transition: non-collinear AFM \rightarrow FM

Magnetoresistance

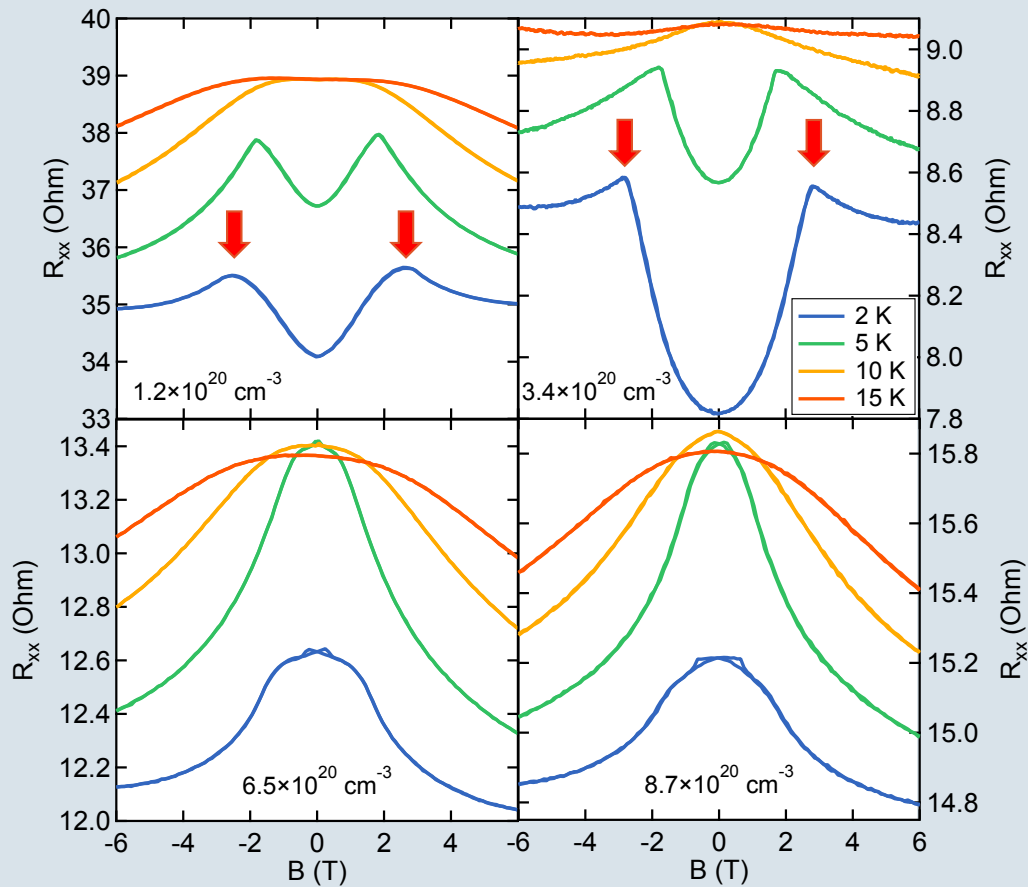


- ☐ Low doping
- ☐ Low-magnetization state
- ☐ Abrupt phase transition with field
- ☐ Metamagnetic transition: non-collinear AFM \rightarrow ?



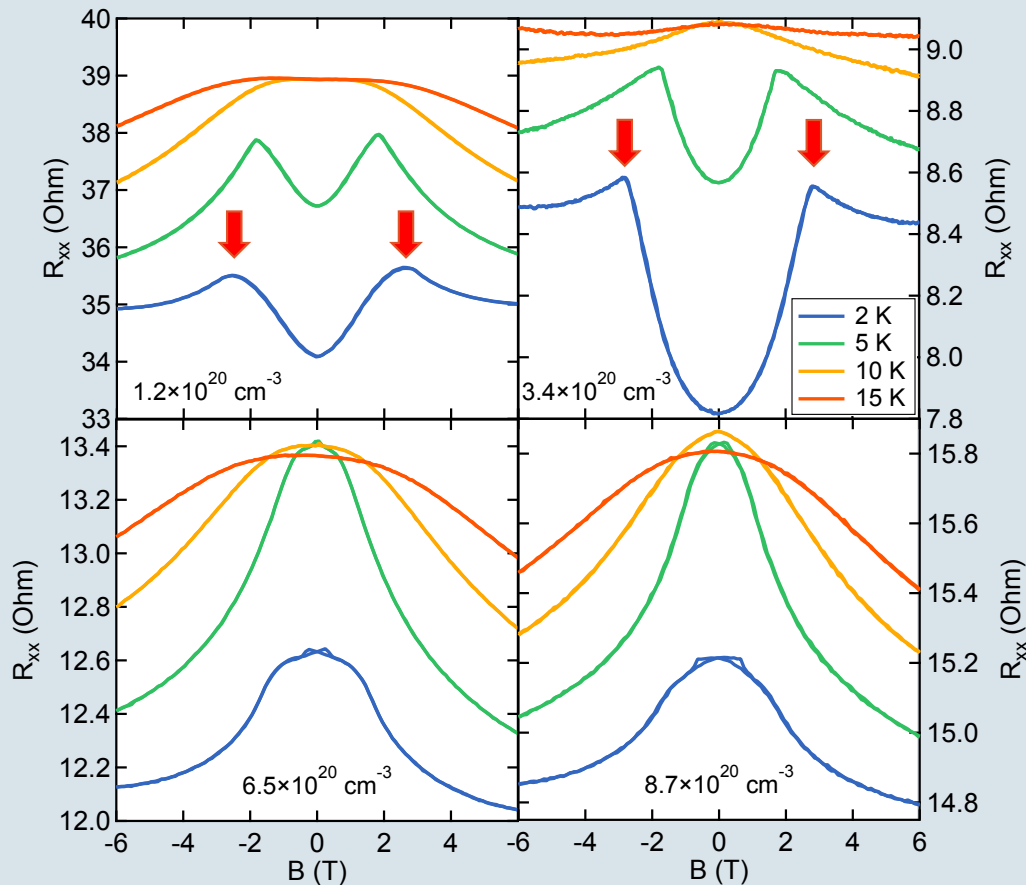
THE vanishes near at the metamagnetic transition

Magnetoresistance



- High doping
- High-magnetization state
- Negative MR: only a gradual change from a non-collinear spin structure to a more collinear one

Magnetoresistance

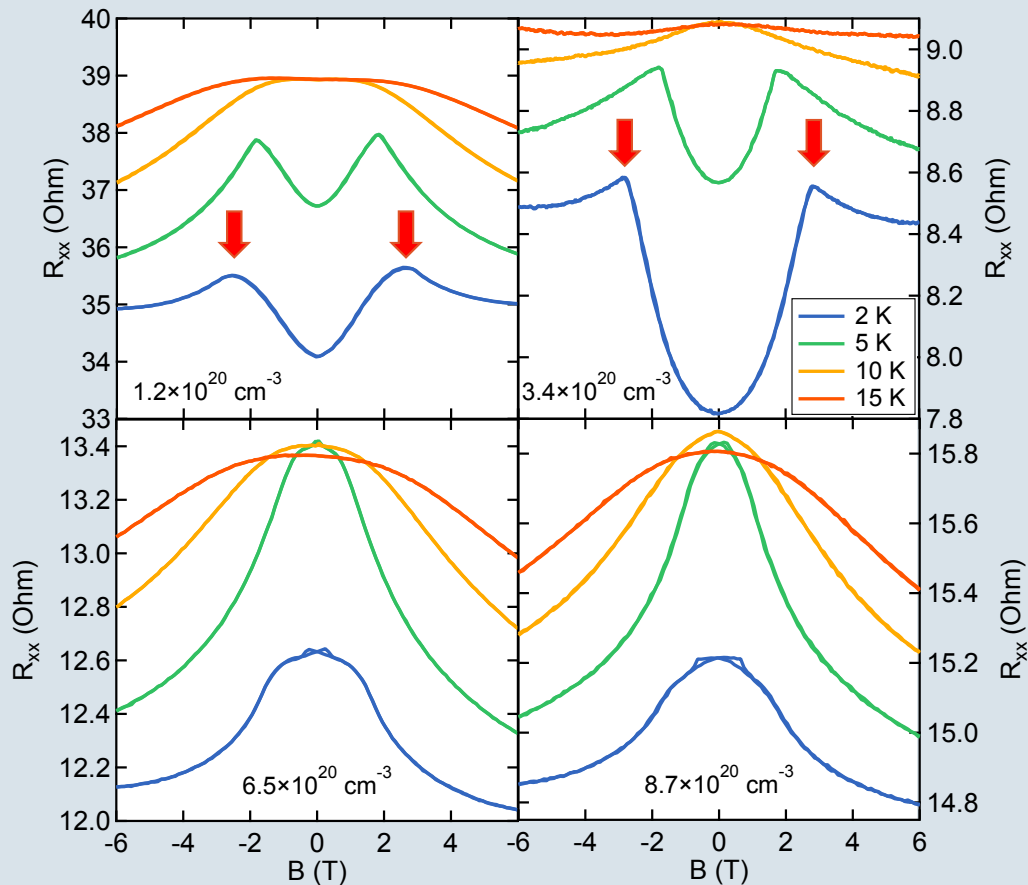


- ☐ Low doping
- ☐ Low-magnetization state
- ☐ Abrupt phase transition with field
- ☐ Metamagnetic transition: non-collinear AFM \rightarrow ?

Presence/absence of phase transition coincides with sign changes in Hall effects, Berry curvature, and Fermi level crossing the Weyl point.

- ☐ High doping
- ☐ High-magnetization state
- ☐ Negative MR: only a gradual change from a non-collinear spin structure to a more collinear one

Magnetoresistance

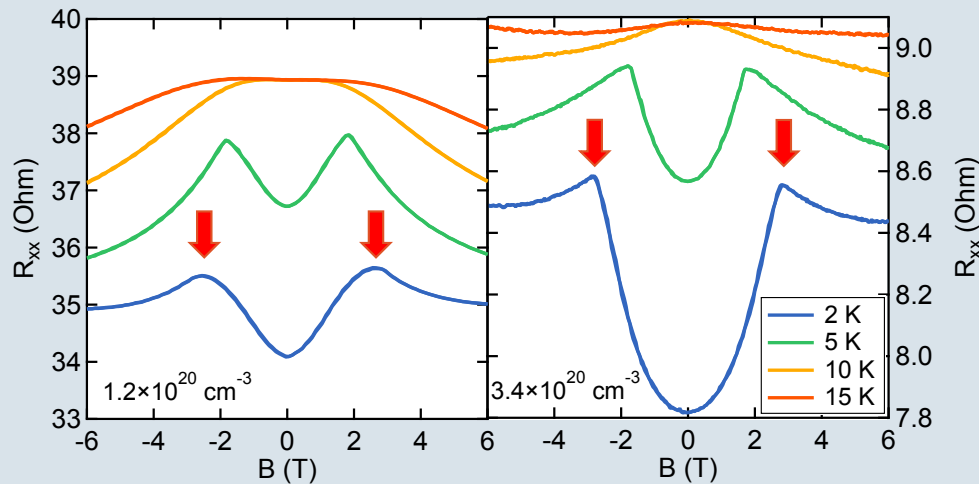


- ☐ Low doping
- ☐ Low-magnetization state
- ☐ Abrupt phase transition with field
- ☐ Metamagnetic transition: non-collinear AFM \rightarrow ?

**Strong connection
between magnetism and
Fermi surface topology.**

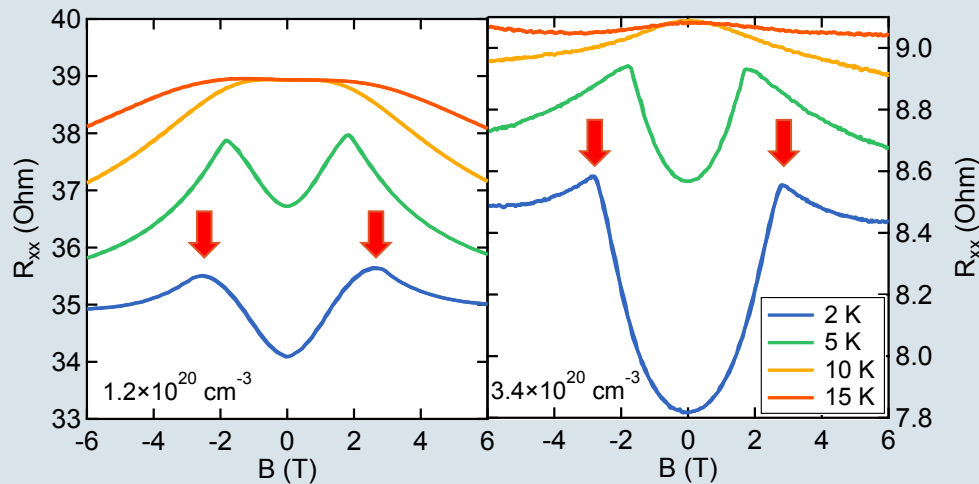
- ☐ High doping
- ☐ High-magnetization state
- ☐ Negative MR: only a gradual change from a non-collinear spin structure to a more collinear one

Metamagnetic Transition

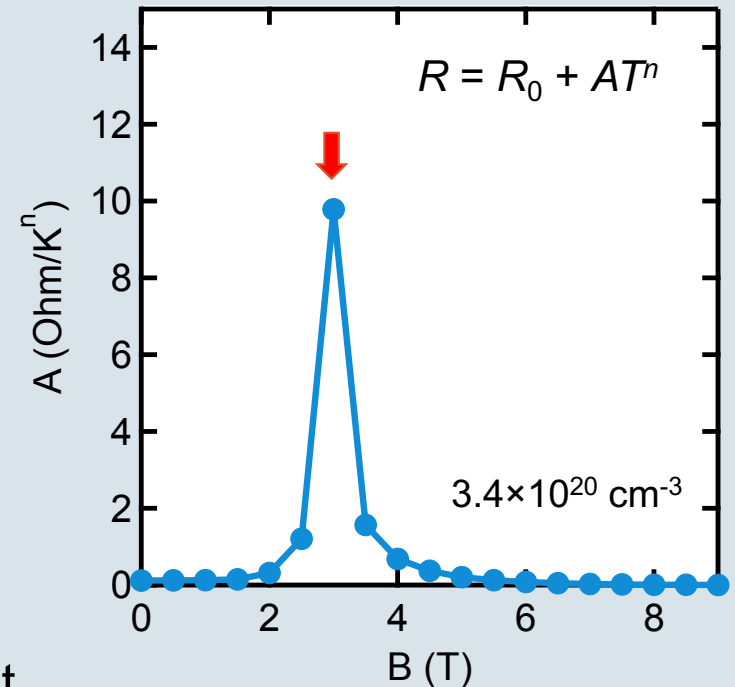


- ❑ Metamagnetic transitions are often associated with singularities in the density of states and Zeeman-splitting induced change in Fermi surface topology (Lifshitz transitions)
- ❑ Presence of Weyl points suggest another possibility: as the magnetic field reorients the spins, it modifies the band topology and thereby magnetotransport

Metamagnetic Transition

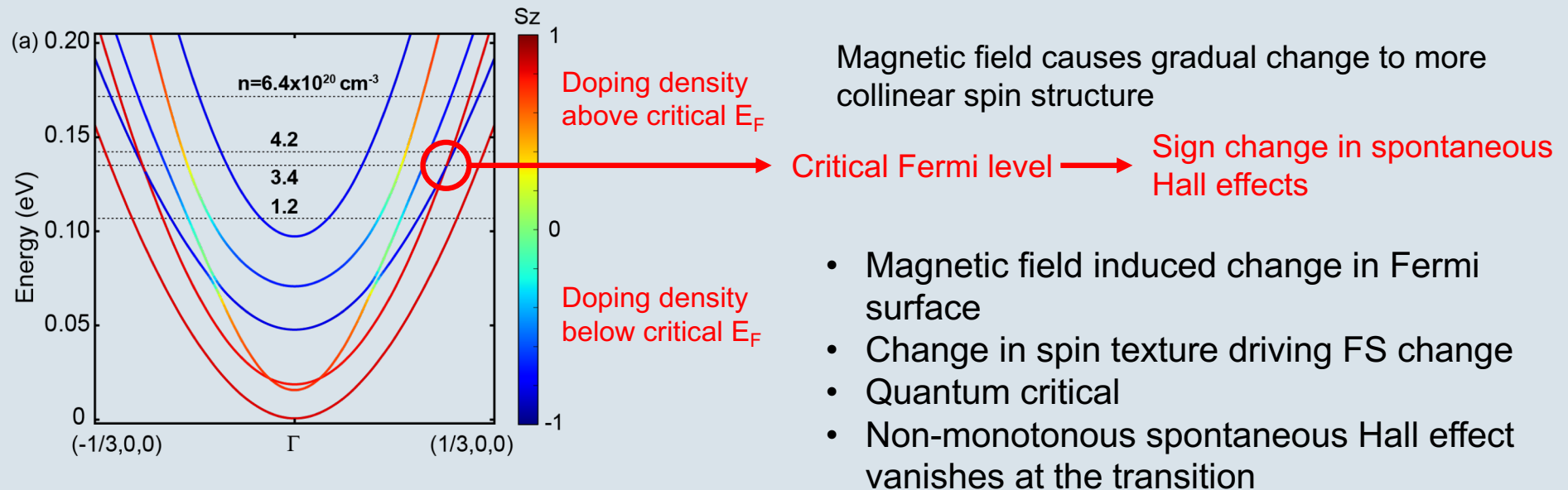


- Peak in A parameter (mass enhancement) near metamagnetic transition suggests magnetic field controlled quantum critical point
- Magnetic field controlled Fermi surface fluctuations/instability

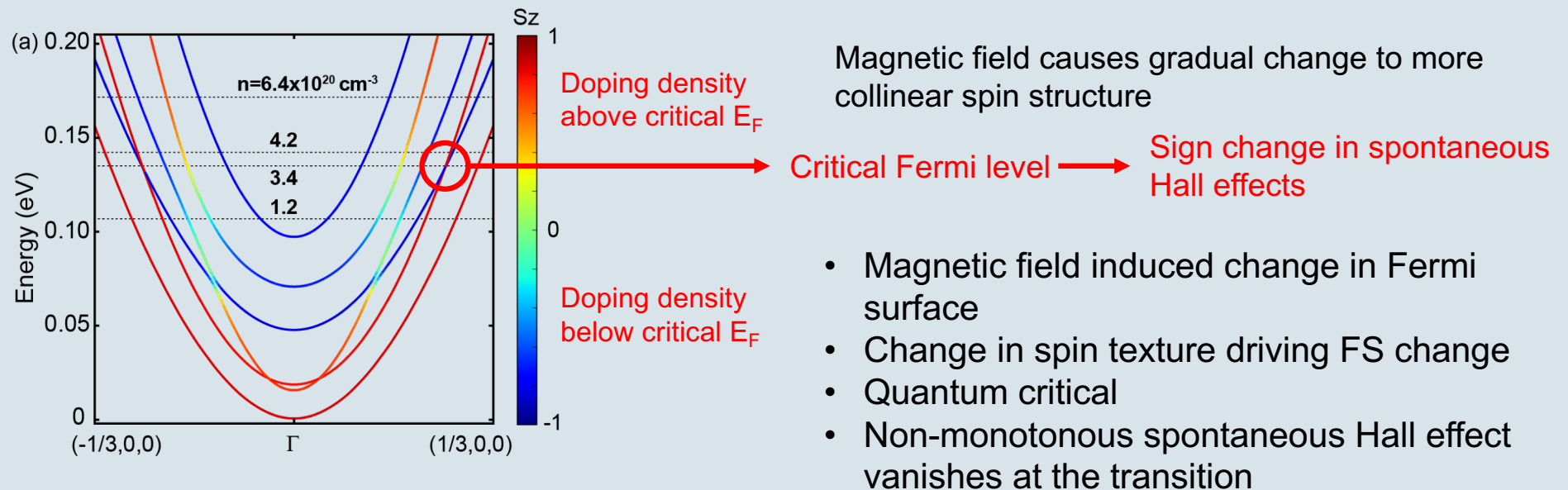


$$\ln \left(\frac{dR(T)}{dT} \right) = \ln(nA) - (n-1) \ln(T)$$

Magnetism and Electronic Structure

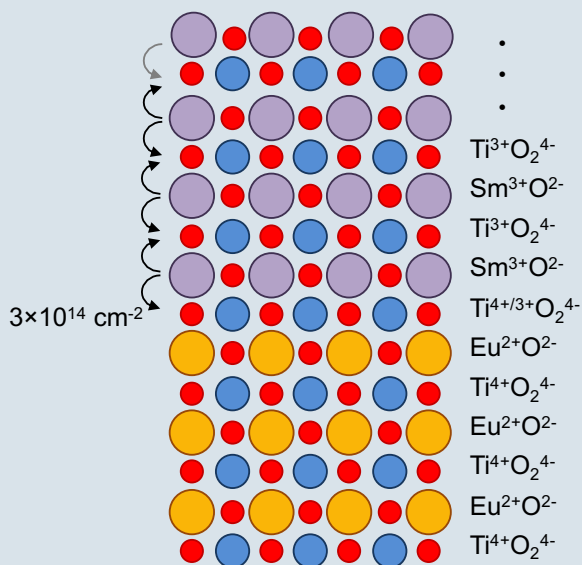
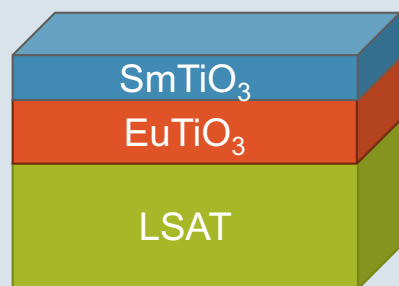


Magnetism and Electronic Structure

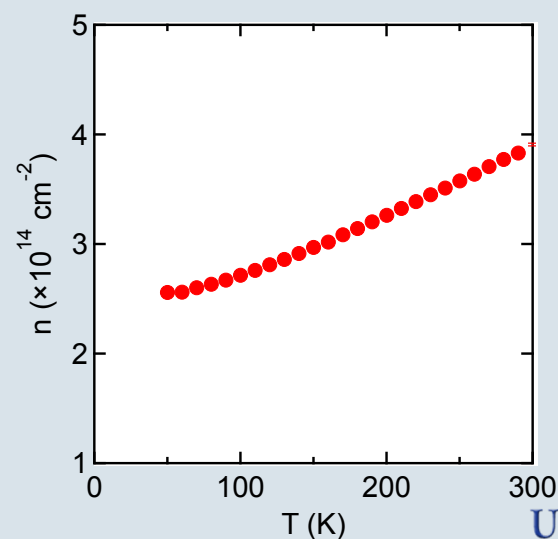
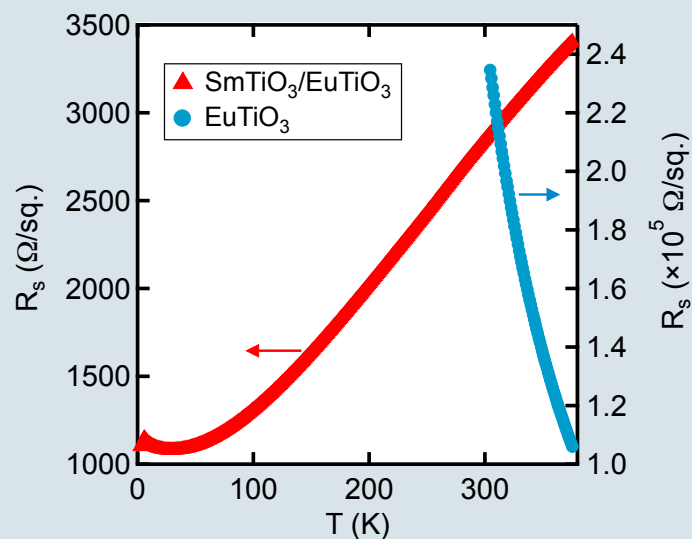


What happens when we modify the electronic band structure, i.e. 3D \rightarrow quasi-2D?

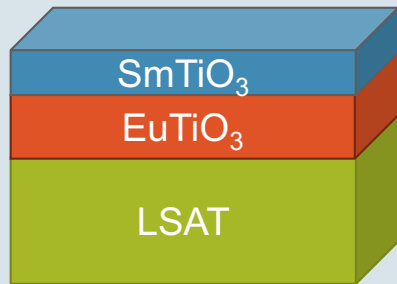
SmTiO₃/EuTiO₃ Interfaces



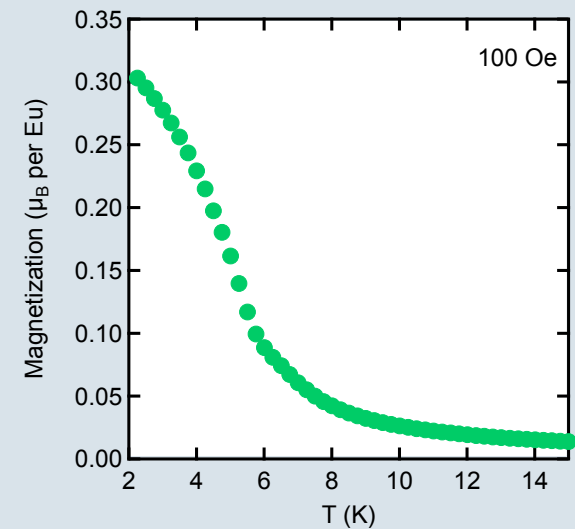
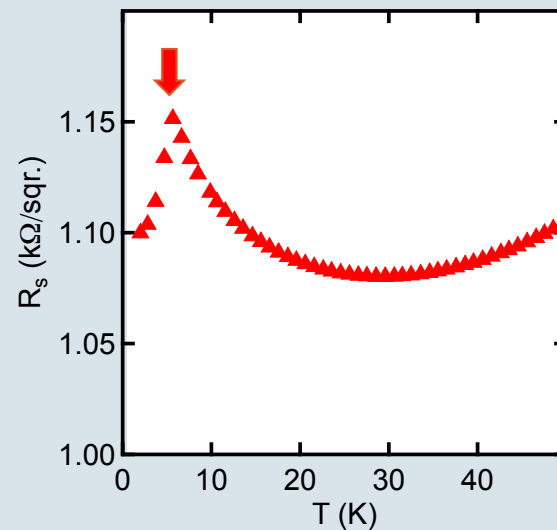
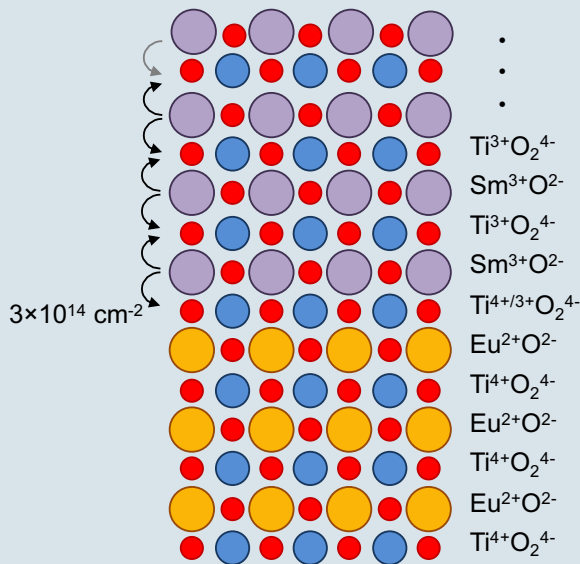
- SmTiO₃/EuTiO₃ interfaces: quasi-two dimensional electron system **in the EuTiO₃**, similar to SmTiO₃/SrTiO₃ interfaces
- Full mobile charge density expected from polar discontinuity: $3.4 \times 10^{14} \text{ cm}^{-2}$
- Charge spreads into the EuTiO₃ layer: “quasi-2D”



SmTiO₃/EuTiO₃ Interfaces



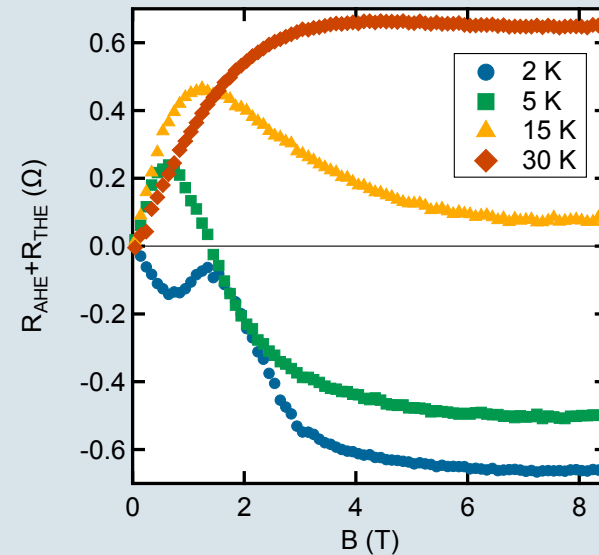
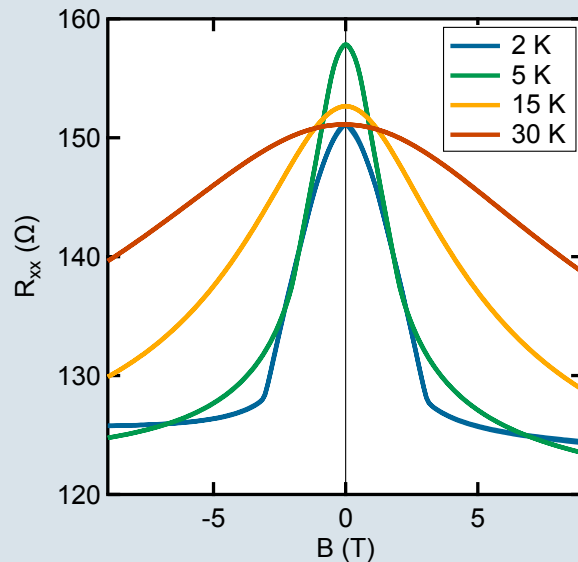
- Similar magnetic transition temperature as bulk EuTiO₃
- Similar magnetization as highly doped EuTiO₃ films
- Mobile charge resides in the EuTiO₃



K. Ahadi, H. Kim, and S. Stemmer, APL Mater. 6, 056102 (2018)

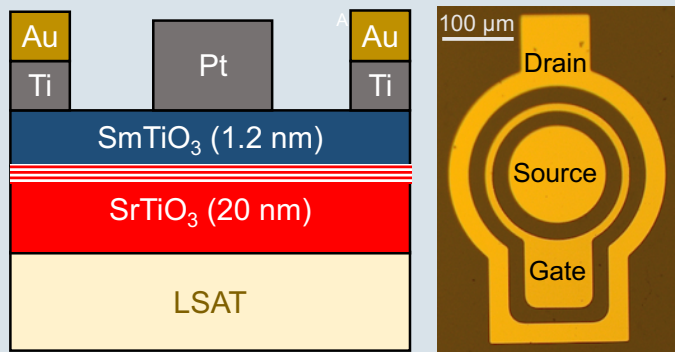
SmTiO₃/EuTiO₃ Interfaces

K. Ahadi, H. Kim,
and S. Stemmer,
APL Mater. 6,
056102 (2018)



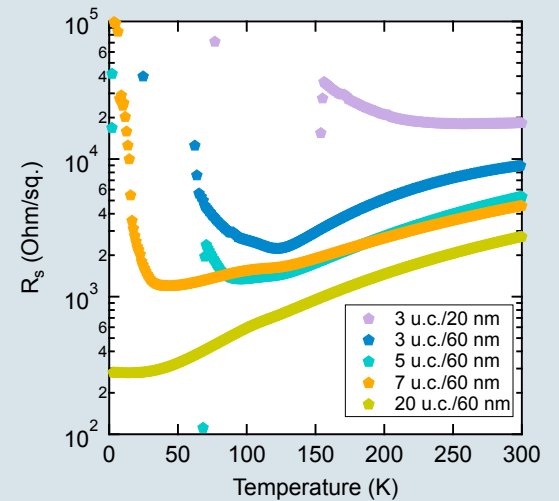
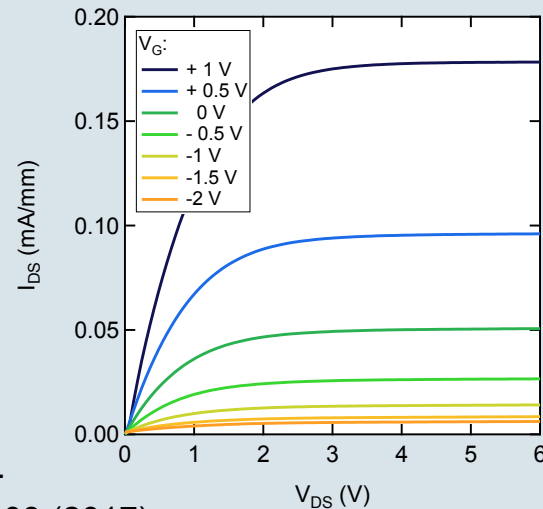
- ❑ AHE changes sign upon lowering the temperature \rightarrow Fermi level closer to points with change in Berry curvature
- ❑ THE appears at higher temperatures than in uniformly doped EuTiO₃, but does not change sign
- ❑ Interface stabilizes non-collinear spin texture to higher temperatures
- ❑ Results consistent with a picture of AHE as a signature of reciprocal space Berry curvature and THE as being due to real space Berry curvature

Outlook



K. Ahadi, et al., Appl. Phys. Lett. 110, 062104 (2017).

K. Ahadi and S. Stemmer, Phys. Rev. Lett. 118, 236803 (2017).



- ☐ Interfaces with thin SmTiO_3 gate dielectrics can be gated (fully depleted)
- ☐ Novel metal-insulator transition: correlation effects?
- ☐ Electron correlations and topology?

Summary

- ❑ Carrier density tuning of spontaneous Hall effects in doped EuTiO_3
- ❑ AHE despite small net magnetization
- ❑ Sign changes in spontaneous Hall effects with carrier density
- ❑ Carrier density control of magnetic field induced transition of the FS and quantum critical points
- ❑ Common origin: band topology (crossings): magnetic topological “semimetal”
- ❑ Results raise interesting questions as to connection between real space magnetic topology and reciprocal space topology
- ❑ **Theory needed**
- ❑ Interesting system for electric field gating

Thank you!